

Hybrid FV-FE methods for modeling flows in fractured porous media

A. Chernyshenko, M. Olshanskii

Marchuk Institute of Numerical Mathematics Russian Academy of Sciences

Multiscale methods and Large-scale Scientific Computing, Moscow,
August 17

Bulk-surface coupled problems

This work develops a hybrid FV–FE method for solving a system of advection–diffusion equations in a bulk domain coupled to advection–diffusion equations on an embedded surface.

Systems of coupled bulk–surface PDEs arise in many engineering and natural science applications:

- the transport and diffusion of solute in fractured porous media
- multiphase fluid dynamics with soluble or insoluble surfactants
- dynamics of biomembranes
- crystal growth
- signaling in biological networks
- etc.

Bulk-surface coupled problems

This work develops a hybrid FV–FE method for solving a system of advection–diffusion equations in a bulk domain coupled to advection–diffusion equations on an embedded surface.

Systems of coupled bulk–surface PDEs arise in many engineering and natural science applications:

- the transport and diffusion of solute in fractured porous media
- multiphase fluid dynamics with soluble or insoluble surfactants
- dynamics of biomembranes
- crystal growth
- signaling in biological networks
- etc.

Numerical treatment of bulk-surface coupled PDEs

Different approaches:

- Fitted FEM: tetrahedral mesh fits the surface + FEM
 - advection-diffusion (Elliott, Ranner, 2013)
 - non-linear reaction-diffusion (Madzvamuse, Chung, 2015,2016)
 - two-phase flow with surfactants (Barrett, Garcke, Nurnberg, 2015)
 - Darcy and transport-diffusion (Alboin, Jaffre, Roberts, Serres, 2002)
- Unfitted FEM: surface cuts through the background tetrahedral mesh
 - cutFEM, Nitsche-XFEM, trace FEM (Burman et al, 2015; Olshanskii, Reusken, Xu, 2014)
 - advection-diffusion (Gross, Olshanskii, Reusken, 2014)
 - two-phase Stokes flow with soluble surfactants (Hansbo, Larson, Zahedi, 2015)
 - time-dependent domains (Hansbo, Larson, Zahedi, 2016)
- FV/FD in bulk domain + FEM (triangulation/extend the PDE off the surface to a narrow band)
- Mixed FEM + Mortar methods (Nordbotten, Boon, Fumagalli, 2017)

Numerical treatment of bulk-surface coupled PDEs

Different approaches:

- Fitted FEM: tetrahedral mesh fits the surface + FEM
 - advection-diffusion (Elliott, Ranner, 2013)
 - non-linear reaction-diffusion (Madzvamuse, Chung, 2015,2016)
 - two-phase flow with surfactants (Barrett, Garcke, Nurnberg, 2015)
 - Darcy and transport-diffusion (Alboin, Jaffre, Roberts, Serres, 2002)
- Unfitted FEM: surface cuts through the background tetrahedral mesh
 - cutFEM, Nitsche-XFEM, trace FEM (Burman et al, 2015; Olshanskii, Reusken, Xu, 2014)
 - advection-diffusion (Gross, Olshanskii, Reusken, 2014)
 - two-phase Stokes flow with soluble surfactants (Hansbo, Larson, Zahedi, 2015)
 - time-dependent domains (Hansbo, Larson, Zahedi, 2016)
- FV/FD in bulk domain + FEM (triangulation/extend the PDE off the surface to a narrow band)
- Mixed FEM + Mortar methods (Nordbotten, Boon, Fumagalli, 2017)

Numerical treatment of bulk-surface coupled PDEs

Different approaches:

- Fitted FEM: tetrahedral mesh fits the surface + FEM
 - advection-diffusion (Elliott, Ranner, 2013)
 - non-linear reaction-diffusion (Madzvamuse, Chung, 2015,2016)
 - two-phase flow with surfactants (Barrett, Garcke,Nurnberg, 2015)
 - Darcy and transport-diffusion (Alboin, Jaffre, Roberts, Serres, 2002)
- Unfitted FEM: surface cuts through the background tetrahedral mesh
 - cutFEM, Nitsche-XFEM, trace FEM (Burman et al, 2015; Olshanskii, Reusken, Xu, 2014)
 - advection-diffusion (Gross, Olshanskii, Reusken, 2014)
 - two-phase Stokes flow with soluble surfactants (Hansbo, Larson, Zahedi, 2015)
 - time-dependent domains (Hansbo, Larson, Zahedi, 2016)
- FV/FD in bulk domain + FEM (triangulation/extend the PDE off the surface to a narrow band)
- Mixed FEM + Mortar methods (Nordbotten, Boon, Fumagalli, 2017)

Numerical treatment of bulk-surface coupled PDEs

Different approaches:

- Fitted FEM: tetrahedral mesh fits the surface + FEM
 - advection-diffusion (Elliott, Ranner, 2013)
 - non-linear reaction-diffusion (Madzvamuse, Chung, 2015,2016)
 - two-phase flow with surfactants (Barrett, Garcke, Nurnberg, 2015)
 - Darcy and transport-diffusion (Alboin, Jaffre, Roberts, Serres, 2002)
- Unfitted FEM: surface cuts through the background tetrahedral mesh
 - cutFEM, Nitsche-XFEM, trace FEM (Burman et al, 2015; Olshanskii, Reusken, Xu, 2014)
 - advection-diffusion (Gross, Olshanskii, Reusken, 2014)
 - two-phase Stokes flow with soluble surfactants (Hansbo, Larson, Zahedi, 2015)
 - time-dependent domains (Hansbo, Larson, Zahedi, 2016)
- FV/FD in bulk domain + FEM (triangulation/extend the PDE off the surface to a narrow band)
- Mixed FEM + Mortar methods (Nordbotten, Boon, Fumagalli, 2017)

Numerical treatment of bulk-surface coupled PDEs

Goals:

- 1 to allow the surface to overlap with the background mesh in an arbitrary way
- 2 avoid building regular surface triangulation
- 3 do not use any extension of the surface PDE to the bulk domain

Approaches:

- octree meshes with cut-cells
- MPFA - nonlinear FV method for polyhedral meshes (Lipnikov, Svyatskiy, Vassilevski, 2012; Chernyshenko, Vassilevski, 2014)
- trace FEM on octree meshes (Chernyshenko, Olshanskii, 2015)

Also receive:

- surface parametrization is not required
- only degrees of freedom from the cells cut by the surface are active

Numerical treatment of bulk-surface coupled PDEs

Goals:

- 1 to allow the surface to overlap with the background mesh in an arbitrary way
- 2 avoid building regular surface triangulation
- 3 do not use any extension of the surface PDE to the bulk domain

Approaches:

- octree meshes with cut-cells
- MPFA - nonlinear FV method for polyhedral meshes (Lipnikov, Svyatskiy, Vassilevski, 2012; Chernyshenko, Vassilevski, 2014)
- trace FEM on octree meshes (Chernyshenko, Olshanskii, 2015)

Also receive:

- surface parametrization is not required
- only degrees of freedom from the cells cut by the surface are active

Numerical treatment of bulk-surface coupled PDEs

Goals:

- 1 to allow the surface to overlap with the background mesh in an arbitrary way
- 2 avoid building regular surface triangulation
- 3 do not use any extension of the surface PDE to the bulk domain

Approaches:

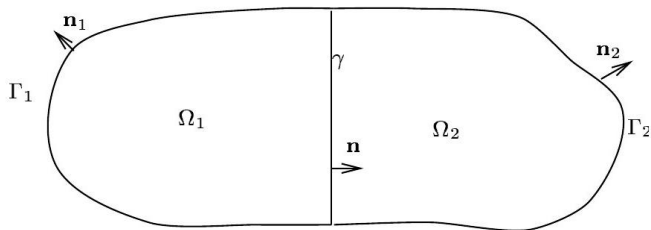
- octree meshes with cut-cells
- MPFA - nonlinear FV method for polyhedral meshes (Lipnikov, Svyatskiy, Vassilevski, 2012; Chernyshenko, Vassilevski, 2014)
- trace FEM on octree meshes (Chernyshenko, Olshanskii, 2015)

Also receive:

- surface parametrization is not required
- only degrees of freedom from the cells cut by the surface are active

Mathematical model

Assume the given bulk domain $\Omega \in \mathbb{R}^3$ and a piecewise smooth surface $\Gamma \subset \Omega$, assume $\partial\Gamma \subset \partial\Omega$. u_i - volume concentration Ω_i , v - surface concentration along Γ .



C. Alboin, J. Jaffre, et.al, *Modeling fractures as interfaces for flow and transport in porous media*, 2001

Mathematical model

Assume the given bulk domain $\Omega \in \mathbb{R}^3$ and a piecewise smooth surface $\Gamma \subset \Omega$, assume $\partial\Gamma \subset \partial\Omega$. u_i - volume concentration Ω_i , v - surface concentration along Γ .

Equations in subdomains:

$$\begin{aligned} \phi_i \frac{\partial u_i}{\partial t} + \operatorname{div}(\mathbf{w}_i u_i) - D_i \Delta u_i &= f_i \quad \text{in } \Omega_i, \\ u_i &= v \quad \text{on } \partial\Omega_i \cap \Gamma \end{aligned} \quad (1)$$

On the surface Γ :

$$\phi_\Gamma \frac{\partial v}{\partial t} + \operatorname{div}_\Gamma(\mathbf{w}_\Gamma v) - d D_\Gamma \Delta_\Gamma v = F_\Gamma(u) + f_\Gamma \quad \text{on } \Gamma, \quad (2)$$

where $\nabla_\Gamma g = \nabla g - \nabla g \cdot \mathbf{n}_\Gamma \mathbf{n}_\Gamma$, $f_\Gamma \in L^2(\Gamma)$, $\mathbf{w}_\Gamma \in H^{1,\infty}(\Gamma)$, $\phi_i > 0, \phi_\Gamma > 0$
- porosity, d - fracture width

Mathematical model

Assume the given bulk domain $\Omega \in \mathbb{R}^3$ and a piecewise smooth surface $\Gamma \subset \Omega$, assume $\partial\Gamma \subset \partial\Omega$. u_i - volume concentration Ω_i , v - surface concentration along Γ .

Equations in subdomains:

$$\begin{aligned} \phi_i \frac{\partial u_i}{\partial t} + \operatorname{div}(\mathbf{w}_i u_i) - D_i \Delta u_i &= f_i \quad \text{in } \Omega_i, \\ u_i &= v \quad \text{on } \partial\Omega_i \cap \Gamma \end{aligned} \quad (1)$$

On the surface Γ :

$$\phi_\Gamma \frac{\partial v}{\partial t} + \operatorname{div}_\Gamma(\mathbf{w}_\Gamma v) - d D_\Gamma \Delta_\Gamma v = F_\Gamma(u) + f_\Gamma \quad \text{on } \Gamma, \quad (2)$$

where $\nabla_\Gamma g = \nabla g - \nabla g \cdot \mathbf{n}_\Gamma \mathbf{n}_\Gamma$, $f_\Gamma \in L^2(\Gamma)$, $\mathbf{w}_\Gamma \in H^{1,\infty}(\Gamma)$, $\phi_i > 0, \phi_\Gamma > 0$
- porosity, d - fracture width

Mathematical model

The total surface flux $F_\Gamma(u)$ represents the contribution of the bulk to the solute transport in the fracture. The mass balance at Γ leads to the equation

$$F_\Gamma(u) = [-D\mathbf{n}_\Gamma \cdot \nabla u + (\mathbf{n}_\Gamma \cdot \mathbf{w})u]_\Gamma, \quad (3)$$

If Γ is piecewise smooth, then we need further conditions on the edges \mathcal{E} .

$$\sum_{j=1}^M \mathbf{w}_{\Gamma,j} \cdot \mathbf{n}_{\partial\Gamma,j} = 0, \quad \sum_{j=1}^M d_j (D_{\Gamma,j} \mathbf{n}_{\partial\Gamma,j}) \cdot \nabla_\Gamma v_j = 0 \quad \text{on } \mathcal{E}. \quad (4)$$

Initial and boundary conditions:

$$\left\{ \begin{array}{l} D_i \mathbf{n}_{\partial\Omega} \cdot \nabla u = 0 \quad \text{on } \partial\Omega_N, \\ u = u_D \quad \text{on } \partial\Omega_D, \\ u|_{t=0} = u_0 \quad \text{in } \Omega, \end{array} \right. \quad \left\{ \begin{array}{l} D_\Gamma \mathbf{n}_{\partial\Gamma} \cdot \nabla_\Gamma v = 0 \quad \text{on } \partial\Gamma_N, \\ v = v_D \quad \text{on } \partial\Gamma_D, \\ v|_{t=0} = v_0 \quad \text{on } \Gamma. \end{array} \right. \quad (5)$$

Mathematical model

The total surface flux $F_\Gamma(u)$ represents the contribution of the bulk to the solute transport in the fracture. The mass balance at Γ leads to the equation

$$F_\Gamma(u) = [-D\mathbf{n}_\Gamma \cdot \nabla u + (\mathbf{n}_\Gamma \cdot \mathbf{w})u]_\Gamma, \quad (3)$$

If Γ is piecewise smooth, then we need further conditions on the edges \mathcal{E} .

$$\sum_{j=1}^M \mathbf{w}_{\Gamma,j} \cdot \mathbf{n}_{\partial\Gamma,j} = 0, \quad \sum_{j=1}^M d_j (D_{\Gamma,j} \mathbf{n}_{\partial\Gamma,j}) \cdot \nabla_\Gamma v_j = 0 \quad \text{on } \mathcal{E}. \quad (4)$$

Initial and boundary conditions:

$$\left\{ \begin{array}{l} D_i \mathbf{n}_{\partial\Omega} \cdot \nabla u = 0 \quad \text{on } \partial\Omega_N, \\ u = u_D \quad \text{on } \partial\Omega_D, \\ u|_{t=0} = u_0 \quad \text{in } \Omega, \end{array} \right. \quad \left\{ \begin{array}{l} D_\Gamma \mathbf{n}_{\partial\Gamma} \cdot \nabla_\Gamma v = 0 \quad \text{on } \partial\Gamma_N, \\ v = v_D \quad \text{on } \partial\Gamma_D, \\ v|_{t=0} = v_0 \quad \text{on } \Gamma. \end{array} \right. \quad (5)$$

Mathematical model

The total surface flux $F_\Gamma(u)$ represents the contribution of the bulk to the solute transport in the fracture. The mass balance at Γ leads to the equation

$$F_\Gamma(u) = [-D\mathbf{n}_\Gamma \cdot \nabla u + (\mathbf{n}_\Gamma \cdot \mathbf{w})u]_\Gamma, \quad (3)$$

If Γ is piecewise smooth, then we need further conditions on the edges \mathcal{E} .

$$\sum_{j=1}^M \mathbf{w}_{\Gamma,j} \cdot \mathbf{n}_{\partial\Gamma,j} = 0, \quad \sum_{j=1}^M d_j (D_{\Gamma,j} \mathbf{n}_{\partial\Gamma,j}) \cdot \nabla_\Gamma v_j = 0 \quad \text{on } \mathcal{E}. \quad (4)$$

Initial and boundary conditions:

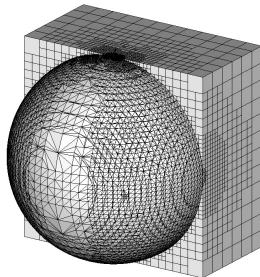
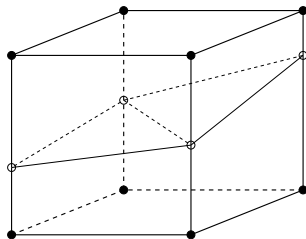
$$\left\{ \begin{array}{l} D_i \mathbf{n}_{\partial\Omega} \cdot \nabla u = 0 \quad \text{on } \partial\Omega_N, \\ u = u_D \quad \text{on } \partial\Omega_D, \\ u|_{t=0} = u_0 \quad \text{in } \Omega, \end{array} \right. \quad \left\{ \begin{array}{l} D_\Gamma \mathbf{n}_{\partial\Gamma} \cdot \nabla_\Gamma v = 0 \quad \text{on } \partial\Gamma_N, \\ v = v_D \quad \text{on } \partial\Gamma_D, \\ v|_{t=0} = v_0 \quad \text{on } \Gamma. \end{array} \right. \quad (5)$$

Approaches. Meshes

Octree meshes with cut-cells (Chernyshenko A., 2013) - modification of Multimaterial marching cubes (Wu Z., Sullivan J.M., 2003) and Cubical marching squares (C.-C. Ho, F.-C.Wu, et al, 2005) algorithms

- \mathcal{T}_h - volume octree mesh in Ω
- Γ intersects \mathcal{T}_h in arbitrary way.
- Γ_h - irregular triangulation of Γ (*2nd order) is used for integration only. Divides Ω into $\Omega_{i,h}$.

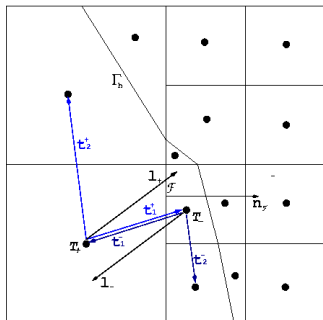
Finally we get a polyhedral mesh (hexahedra + polyhedra)



Approaches. Bulk: Monotone finite volume method

Monotone FV method for the advection-diffusion equations

- 2D: K. Lipnikov, D. Svyatskiy, Yu. Vassilevski, 2012
- 3D advection: Nikitin K., Vassilevski Yu. , 2010
- 3D diffusion: Chernyshenko A., Vassilevski Yu. , 2014



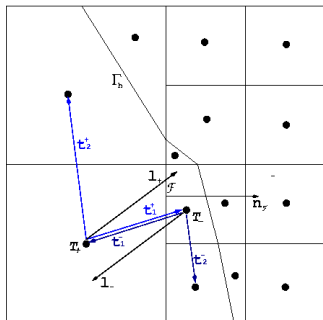
- nonlinear
- compact stencil
- monotone (DMP)
- *2nd order in concentrations

A.Y.Chernyshenko, Yu.V.Vassilevski, *Finite volume scheme with the discrete maximum principle for diffusion equations on polyhedral meshes.*// FVCA7, 2014.

Approaches. Bulk: Monotone finite volume method

Monotone FV method for the advection-diffusion equations

- 2D: K. Lipnikov, D. Svyatskiy, Yu. Vassilevski, 2012
- 3D advection: Nikitin K., Vassilevski Yu. , 2010
- 3D diffusion: Chernyshenko A., Vassilevski Yu. , 2014

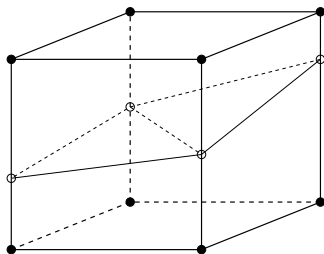


- Diffusive flux $\mathbf{q} = -D\nabla u$
- $D\nabla u \cdot \mathbf{n}_f = \nabla u \cdot (D\mathbf{n}_f) = \nabla u \cdot \mathbf{l}_F$
- $\mathbf{l}_F = \alpha_+ \mathbf{t}_1^+ + \beta_+ \mathbf{t}_2^+$
- $q_+ = A_+ (u_+ - u_-) + B_+ (u_+ - u_{+,2})$
 $q_- = A_- (u_- - u_+) + B_- (u_- - u_{-,2})$

A.Y.Chernyshenko, Yu.V.Vassilevski, *Finite volume scheme with the discrete maximum principle for diffusion equations on polyhedral meshes*,// FVCA7, 2014.

Approaches. Surface: The trace finite element method.

- \mathcal{T}_h - octree cubic mesh covering Ω
- $\Gamma_h = \bigcup_{K \in \mathcal{F}_h} K$ - approximation of Γ . For any $K \in \mathcal{F}_h$ there is only *one* cell $T_K \in \mathcal{T}_h$ such that $K \subset T_K$
- $\tilde{\Gamma}_h := \{\mathbf{x} \in \Omega : \phi_h(\mathbf{x}) = 0\}$, where ϕ_h - interpolant to ϕ - level set function of Γ . Finding ϕ_h , we recover Γ_h by Cubical marching squares.



Approaches. Surface: The trace finite element method.

- V_h - volumetric finite element space .
 $V_h := \{v_h \in C(\Omega) \mid v_h|_S \in Q_1 \ \forall S \in \mathcal{T}_h\}$, where Q_1 -all piecewise trilinear continuous functions with respect to the bulk octree mesh \mathcal{T}_h .
- V_h^Γ - the space of traces on Γ_h of all piecewise trilinear continuous functions with respect to the outer triangulation \mathcal{T}_h
 $V_h^\Gamma := \{\psi_h \in H^1(\Gamma_h) \mid \exists v_h \in V_h \text{ such that } \psi_h = v_h|_{\Gamma_h}\}$.
- FE-discretization of (2): Find $v_h \in V_h^\Gamma$ such that

$$\int_{\Gamma_h} \left(\phi_{\Gamma,h} \frac{\partial v_h}{\partial t} w_h + d_h D_{\Gamma,h} \nabla_{\Gamma_h} v_h \cdot \nabla_{\Gamma_h} w_h + (\mathbf{w}_h \cdot \nabla_{\Gamma_h} v_h) w_h \right) + (\operatorname{div}_{\Gamma_h} \mathbf{w}_h) u_h v_h \, ds_h = \int_{\Gamma_h} (F_{\Gamma,h}(u_h) + f_{\Gamma,h}) w_h \, ds_h \quad (6)$$

for all $w_h \in V_h^\Gamma$

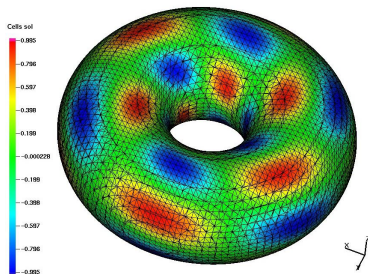
The trace finite element method. Variants

- Uniform refinement or local refinement depending on error indicators and surface curvatures
- Full gradient ∇ or tangential gradient ∇_Γ ($\nabla_\Gamma u = \nabla u^{ext}$).
- SUPG stabilized method

Numerical analysis and error estimates can be found in

A.Y.Chernyshenko, M.A. Olshanski, *An adaptive octree finite element method for PDEs posed on surfaces* //CMAME, 2015.

The trace finite element method. Numerical experiments



$$-\Delta_{\Gamma} u = f \quad \text{on } \Gamma,$$

where

$$\Gamma = \{\mathbf{x} \in \Omega \mid r^2 = x_3^2 + (\sqrt{x_1^2 + x_2^2} - R)^2\},$$

$$\Gamma \subset \Omega = (-2, 2)^3,$$

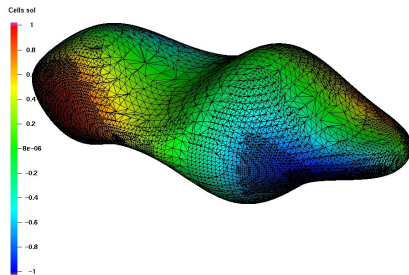
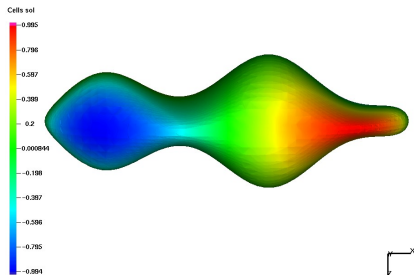
$$u(\mathbf{x}) = \sin(3\phi) \cos(3\theta + \phi)$$

#d.o.f.	surface gradient variant				full gradient variant			
	L^2 -norm	rate	L^∞ -norm	rate	L^2 -norm	rate	L^∞ -norm	rate
297	2.415e-1		7.898e-1		3.199e-1		8.861e-1	
1289	5.047e-2	2.26	2.123e-1	1.90	8.406e-2	1.93	4.437e-1	0.98
5001	1.049e-2	2.27	5.400e-2	1.98	1.956e-2	2.10	8.562e-2	2.37
20073	2.367e-3	2.15	1.341e-2	2.01	4.979e-3	1.97	2.184e-2	1.97

The trace finite element method. Numerical experiments

$$\phi(\mathbf{x}) = \frac{1}{4}x_1^2 + x_2^2 + \frac{4x_3^2}{(1 + \frac{1}{2}\sin(\pi x_1))^2} - 1.$$

$\mathbf{w}_\Gamma = 0$, $D_\Gamma = 1$, $u = x_1x_2$ on Γ .

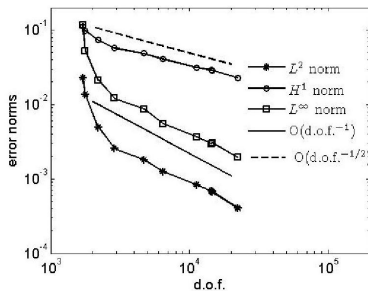
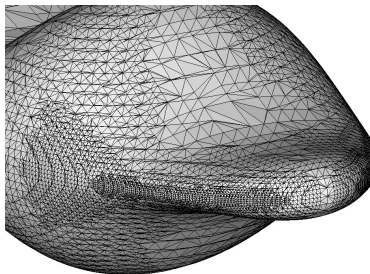


G. Dziuk and C. M. Elliott, *Finite element methods for surface PDEs*, Acta Numerica (2013).

The trace finite element method. Numerical experiments

$$\phi(\mathbf{x}) = \frac{1}{4}x_1^2 + x_2^2 + \frac{4x_3^2}{\left(1 + \frac{1}{2}\sin(\pi x_1)\right)^2} - 1.$$

$\mathbf{w}_\Gamma = 0$, $D_\Gamma = 1$, $u = x_1x_2$ on Γ .



Return to bulk-surface coupled problem

Coupling between discrete bulk and surface equations

- Boundary condition $u_i = v$ on $\partial\Omega_{i,h} \cap \Gamma_h$
- Net flux $F_{\Gamma_h}(u)$ on Γ_h

We assume an implicit time stepping method.

This results in the following system on each time step.

$$\left\{ \begin{array}{ll} \mathcal{L}u := \tilde{\phi}u + \operatorname{div}(\mathbf{w}u - D\nabla)u = \hat{f} & \text{in } \Omega \setminus \Gamma, \\ u = v & \text{on } \Gamma, \\ D\mathbf{n}_{\partial\Omega} \cdot \nabla u = 0 & \text{on } \partial\Omega_N, \quad u = u_D & \text{on } \partial\Omega_D, \\ \mathcal{L}_\Gamma v := \tilde{\phi}_\Gamma v + \operatorname{div}_\Gamma(\mathbf{w}_\Gamma v - dD_\Gamma \nabla_\Gamma v) = F_\Gamma(u) + \hat{f}_\Gamma & \text{on } \Gamma, \\ D_\Gamma \mathbf{n}_{\partial\Gamma} \cdot \nabla v = 0 & \text{on } \partial\Gamma_N, \quad v = v_D & \text{on } \partial\Gamma_D, \end{array} \right.$$

Coupling between discrete bulk and surface equations

We solve the coupled system by the fixed point method. Given u^0, v^0 , the initial guess, we iterate for $k = 0, 1, 2, \dots$ until convergence:

Step 1: Solve for u^{k+1} ,

$$\begin{cases} \mathcal{L}u^{k+1} = \hat{f} \text{ in } \Omega \setminus \Gamma, & u^{k+1} = v^k \text{ on } \Gamma, \\ D\mathbf{n}_{\partial\Omega} \cdot \nabla u^{k+1} = 0 \text{ on } \partial\Omega_N, & u^{k+1} = u_D \text{ on } \partial\Omega_D, \end{cases} \quad (7)$$

Step 2: Solve for v^{aux} and update for v^{k+1} with a relaxation parameter ω ,

$$\begin{cases} \mathcal{L}_\Gamma v^{\text{aux}} = F_\Gamma(u^{k+1}) + \hat{f}_\Gamma \text{ on } \Gamma, \\ D_\Gamma \mathbf{n}_{\partial\Gamma} \cdot \nabla_\Gamma v^{\text{aux}} = 0 \text{ on } \partial\Gamma_N, & v^{\text{aux}} = v_D \text{ on } \partial\Gamma_D \\ v^{k+1} = \omega v^{\text{aux}} + (1 - \omega)v^k, & \omega \in (0, 1], \end{cases} \quad (8)$$

Coupling between discrete bulk and surface equations

We solve the coupled system by the fixed point method. Given u^0, v^0 , the initial guess, we iterate for $k = 0, 1, 2, \dots$ until convergence:

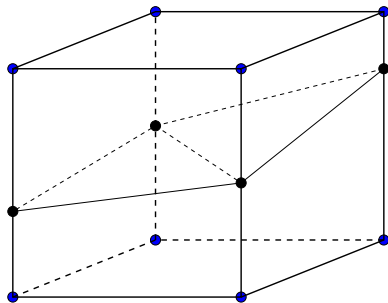
Step 1: Solve for u^{k+1} ,

$$\begin{cases} \mathcal{L}u^{k+1} = \hat{f} \text{ in } \Omega \setminus \Gamma, & u^{k+1} = v^k \text{ on } \Gamma, \\ D\mathbf{n}_{\partial\Omega} \cdot \nabla u^{k+1} = 0 \text{ on } \partial\Omega_N, & u^{k+1} = u_D \text{ on } \partial\Omega_D, \end{cases} \quad (7)$$

Step 2: Solve for v^{aux} and update for v^{k+1} with a relaxation parameter ω ,

$$\begin{cases} \mathcal{L}_\Gamma v^{\text{aux}} = F_\Gamma(u^{k+1}) + \hat{f}_\Gamma \text{ on } \Gamma, \\ D_\Gamma \mathbf{n}_{\partial\Gamma} \cdot \nabla_\Gamma v^{\text{aux}} = 0 \text{ on } \partial\Gamma_N, & v^{\text{aux}} = v_D \text{ on } \partial\Gamma_D \\ v^{k+1} = \omega v^{\text{aux}} + (1 - \omega)v^k, & \omega \in (0, 1], \end{cases} \quad (8)$$

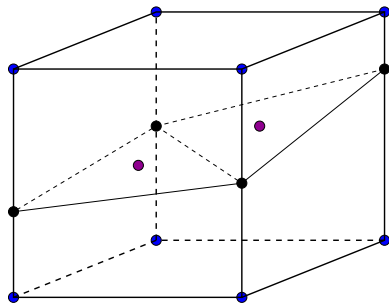
Coupling between discrete bulk and surface equations



● DOFs for trace FEM

- Solve PDE on surface using TraceFEM

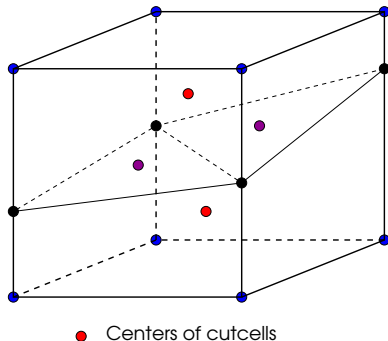
Coupling between discrete bulk and surface equations



● Centers of triangles

- Get solution at centers of triangles

Coupling between discrete bulk and surface equations



- Solve 3D problems in subdomains
- Calculate fluxes through the triangles

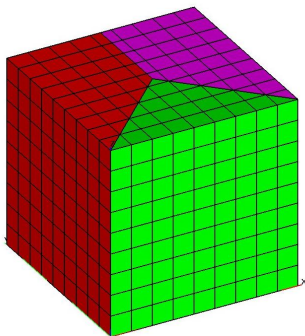
$\Omega = [0, 1]^3$, $\Gamma = \Gamma_{12} \cup \Gamma_{13} \cup \Gamma_{23}$ - divides Ω into three parts
 Consider functions

$$\phi_1 = \begin{cases} 16(y - y_0)^4, & y > y_0 \\ 0, & y \leq y_0 \end{cases}$$

and $\phi_2 = x - y$, $\phi_3 = x + y - 1$.

For the exact solution which is continuous but has derivative jump we take the function

$$\begin{cases} u_1 = \sin(2\pi z) \cdot \phi_2 \cdot \phi_3 & \text{in } \Omega_1 \\ u_2 = \sin(2\pi z) \cdot \phi_1 & \text{in } \Omega_2 \\ u_3 = \sin(2\pi z) 2x \cdot \phi_1 & \text{in } \Omega_3 \end{cases}$$



$\Omega = [0, 1]^3$, $\Gamma = \Gamma_{12} \cup \Gamma_{13} \cup \Gamma_{23}$ - divides Ω into three parts
 Consider functions

$$\phi_1 = \begin{cases} 16(y - y_0)^4, & y > y_0 \\ 0, & y \leq y_0 \end{cases}$$

and $\phi_2 = x - y$, $\phi_3 = x + y - 1$.
 For the exact solution which is continuous but has derivative jump we take the function

$$\begin{cases} u_1 = \sin(2\pi z) \cdot \phi_2 \cdot \phi_3 & \text{in } \Omega_1 \\ u_2 = \sin(2\pi z) \cdot \phi_1 & \text{in } \Omega_2 \\ u_3 = \sin(2\pi z) 2x \cdot \phi_1 & \text{in } \Omega_3 \end{cases}$$

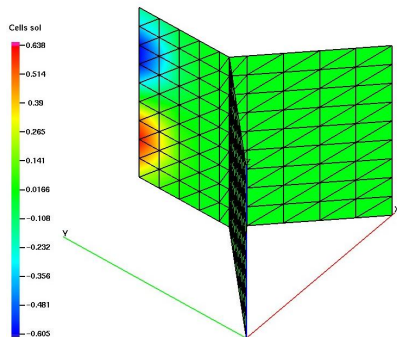


Table: Convergence of numerical solutions

	#d.o.f.	L^2 -norm	rate	H^1 -norm	rate	L^∞ -norm	rate
3D	855	6.374e-3		4.214e-1		3.920e-2	
	7410	1.698e-3	1.84	1.631e-1	1.36	1.276e-2	1.56
	61620	4.235e-4	1.97	6.193e-2	1.39	3.506e-3	1.83
	502440	1.044e-4	2.00	2.348e-2	1.40	1.129e-3	1.62
Surface	232	8.469e-3		2.914e-1		9.280e-3	
	1242	2.003e-3	1.79	1.387e-1	0.92	2.779e-3	1.44
	5662	5.588e-4	1.84	6.874e-2	1.01	1.217e-3	1.09
	24102	1.791e-4	1.64	3.395e-2	1.02	5.181e-4	1.18

Next we rotate the fracture around the line $x = 0.5, z = 0.5$ by angle α .

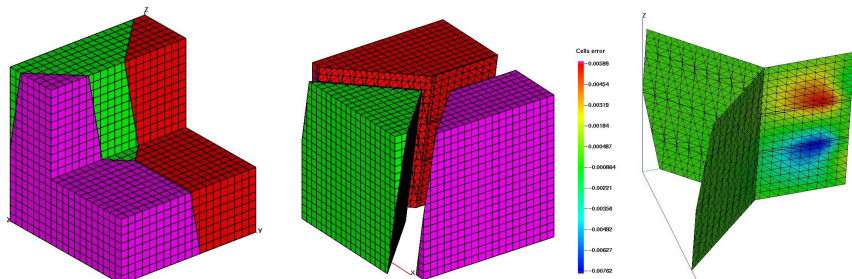


Figure: $\alpha = 20$. 3D mesh, error on rotated fracture (right)

Next we rotate the fracture around the line $x = 0.5, z = 0.5$ by angle α .

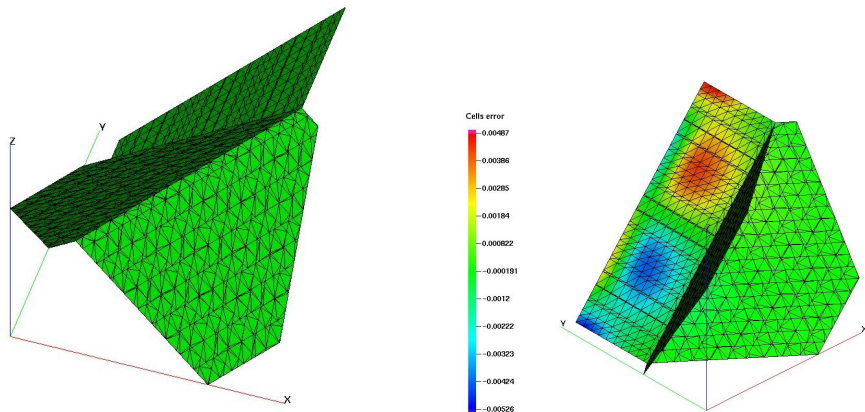


Figure: $\alpha = 40$. Surface mesh, error on rotated fracture (right)

Next we rotate the fracture around the line $x = 0.5, z = 0.5$ by angle α .

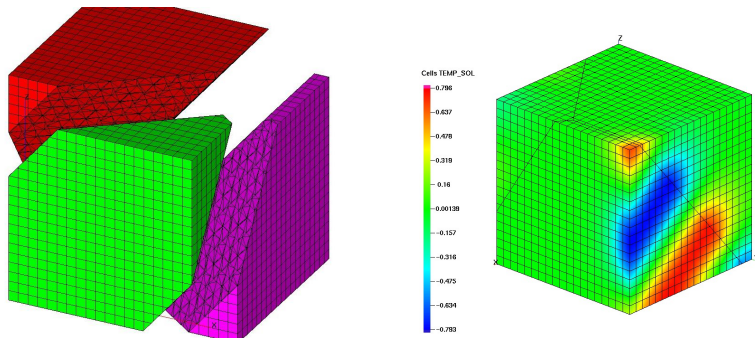


Figure: $\alpha = 40$. 3D mesh, solution on 3D mesh

Table: Convergence of numerical solutions. $\alpha = 20$

	#d.o.f.	L^2 -norm	rate	H^1 -norm	rate	L^∞ -norm	rate
3D	965	6.319e-3		4.208e-1		3.754e-2	
	7872	1.805e-3	1.79	1.661e-1	1.34	1.280e-2	1.55
	63592	5.623e-4	1.80	6.371e-2	1.38	3.411e-3	1.90
	510390	1.602e-4	1.81	2.442e-2	1.39	1.146e-3	1.57
Surface	321	7.792e-3		2.694e-1		2.716e-2	
	1692	2.084e-3	1.59	1.240e-1	1.03	5.400e-3	1.94
	7944	7.019e-4	1.41	6.291e-2	0.97	2.001e-3	1.29
	33272	2.441e-4	1.52	3.173e-2	0.99	7.217e-3	1.47

Table: Convergence of numerical solutions. $\alpha = 40$

	#d.o.f.	L^2 -norm	rate	H^1 -norm	rate	L^∞ -norm	rate
3D	991	5.934e-3		4.080e-1		3.783e-2	
	7996	1.700e-3	1.80	1.621e-1	1.33	1.276e-2	1.56
	64046	4.907e-4	1.80	6.263e-2	1.37	3.515e-3	1.86
	512258	1.503e-4	1.82	2.541e-2	1.39	1.237e-3	1.61
Surface	353	8.167e-3		2.709e-1		2.696e-2	
	1932	2.146e-3	1.66	1.275e-1	1.02	5.566e-3	1.85
	8766	7.115e-4	1.59	6.279e-2	0.97	2.063e-3	1.31
	36676	2.538e-4	1.49	3.121e-2	1.01	7.251e-4	1.51

Boundary conditions: on the $y = 1$ we have $u = 1$ - constant concentration of a contaminant , on others: $\frac{\partial u}{\partial \mathbf{n}_{\partial\Omega}} = 0$

$u(0) = 0$ in Ω and $v(0) = 0$ on Γ .

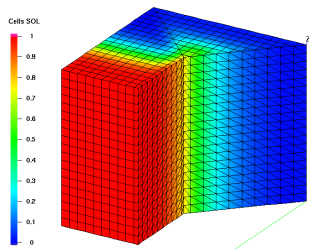
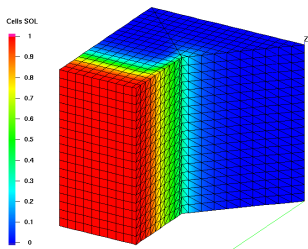
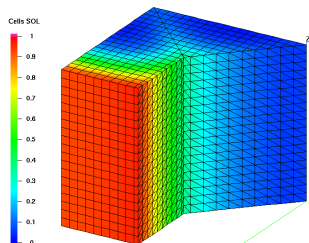
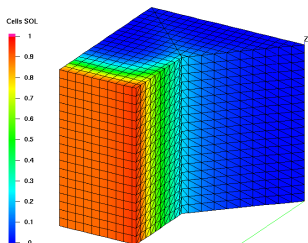
The velocity field: $\mathbf{w} = 2k(0, -1, 0)^T$ in Ω

$$\begin{cases} \mathbf{w}_{\Gamma,1} &= 5k(0, -1, 0), \\ \mathbf{w}_{\Gamma,2} &= 5k(-1/\sqrt{2}, -1/\sqrt{2}, 0), \\ \mathbf{w}_{\Gamma,3} &= 5k(1/\sqrt{2}, -1/\sqrt{2}, 0), \end{cases}$$

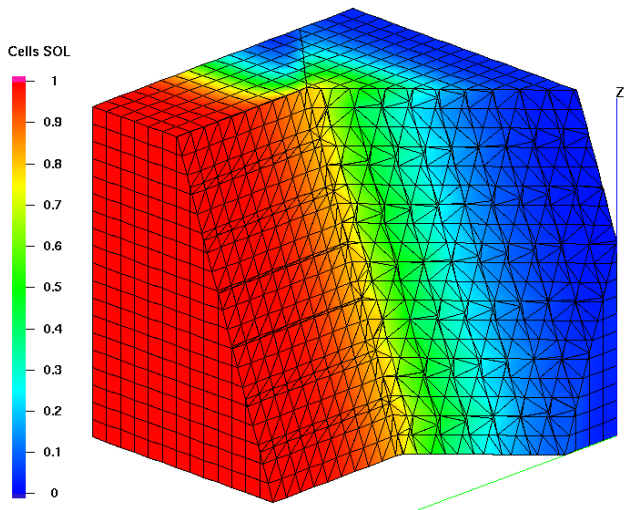
For the case of diffusion domination $k = 1/8$, advection-dominated problem $k = 8$.

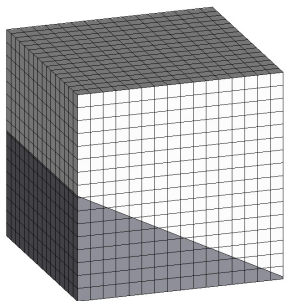
$$D_1 = D_2 = 0.1 I, \quad d = 1, \quad D_{\Gamma} = I, \quad \phi_1 = \phi_2 = \phi_{\Gamma} = 1.$$

Numerical examples. Triple fracture, unsteady problem



Numerical examples. Triple fraction, unsteady problem





$$\Omega = (0, 1)^3, \text{ and}$$

$$\Gamma = \{\mathbf{x} \in \Omega : z + \frac{1}{2}x = 0.51\}.$$

$x = 0$ - inflow

$\mathbf{w}_i = 0$ in Ω_i (no flow in the rock),

$$D_i = 10^{-6}I, D_\Gamma = 10^{-4}I$$

In the fracture the wind is constant $\mathbf{w}_\Gamma = (w_1, 0, w_3)$, $|\mathbf{w}_\Gamma| = 1$, and the contaminant source occupies the part of $\partial\Gamma$,

$$\partial\Gamma_D = \{(0, y, 0.51) : y \in (\frac{1}{4}, \frac{3}{4})\}, v_D = 1 \text{ on } \partial\Gamma_D.$$

This is a bulk–surface variant of a standard test case of numerical solvers for convection–diffusion problems (N.-E. Sun, A.Sun, 2013)

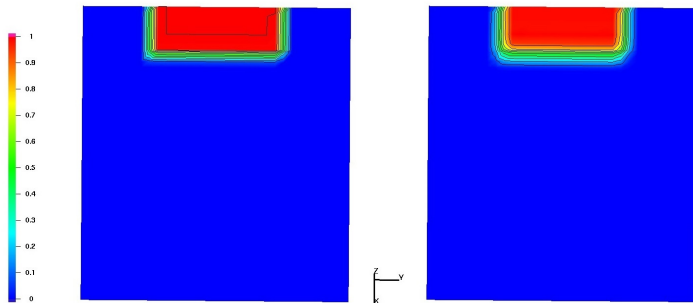


Figure: Reference 2D solution (left) and the fracture component of the computed solution (right) for the contaminant transport along the fracture test case. The solutions is shown for the time $t = 0.17$.

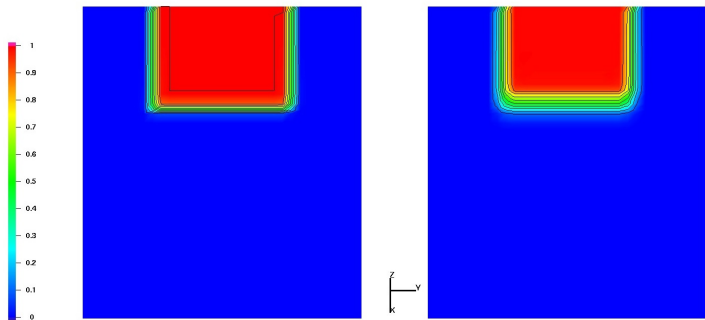


Figure: Reference 2D solution (left) and the fracture component of the computed solution (right) for the contaminant transport along the fracture test case. The solutions is shown for the times $t = 0.34$.

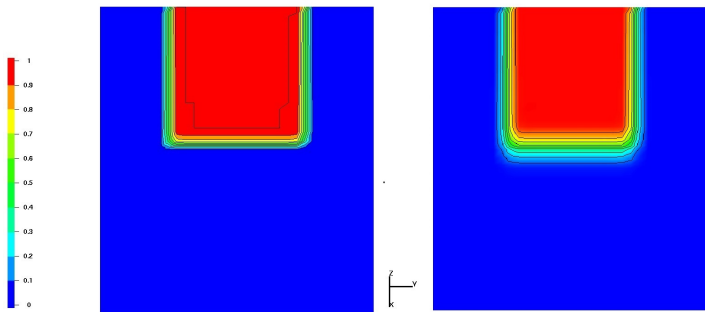


Figure: Reference 2D solution (left) and the fracture component of the computed solution (right) for the contaminant transport along the fracture test case. The solutions is shown for the times $t = 0.5$.

Conclusions

- We studied a hybrid FV-FE method for bulk-surface coupled problems
- The method based on: octree meshes with cut-cells, non-linear FV scheme for polyhedral meshes, trace octree FEM
- The 3D-mesh is unfitted to a surface, the method works for surfaces defined implicitly, parametrization of a surface is not required
- Numerical experiments shows optimal order of convergence

Further work

- Stabilized method for Darcy flow
- Realistic geometries

Chernyshenko A.Yu., Olshanskii M.A., Vassilevski Yu.V, *A hybrid finite volume – finite element method for bulk–surface coupled problems.* // JCP, V. 352, 2018.

Thank You for attention!